

## The Crystal Structure of Trimethylselenonium Iodide, $(\text{CH}_3)_3\text{SeI}$

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(Received 19 July 1965)

Crystals of trimethylselenonium iodide,  $(\text{CH}_3)_3\text{SeI}$  are orthorhombic, with space group  $Pnma$ , cell dimensions  $a=14.078$ ,  $b=8.000$ , and  $c=6.177$  Å (formal e.s.d.  $\sim 0.01\%$ ), and four formula units in the cell. The structure was refined by a three-dimensional full-matrix least-squares procedure. The

structure is built up of pairs of selenium and iodide ions with a linear arrangement of C—Se ··· I,

where C—Se = 1.962 ( $\sigma=0.023$ ) and Se—I = 3.776 ( $\sigma=0.002$ ). The other C—Se distances are 1.946 ( $\sigma=0.016$ ) Å. The mean value of the nearly equal C—Se—C angles is  $98.5^\circ$ .

### Introduction

Both sulfur and selenium form stable 'onium' compounds of the type  $\text{R}_3\text{SX}$  or  $\text{R}_3\text{SeX}$ . The geometry of a sulfonium ion was established through the determination of the crystal structure of trimethylsulfonium iodide by Zuccaro & McCullough (1959). Only an incomplete study on one selenium compound, triphenylselenonium chloride, appears to have been reported in the literature (McCullough & Marsh, 1950*a*). The Se—Cl distance of 3.6 Å indicates an ionic compound, but no other conclusions regarding the structure of the selenium ion could be drawn from that study. It was therefore felt that an investigation of the crystal structure of a salt of this type was desirable. Among the several selenium salts that were readily available, it was found that the trimethylselenonium iodide would be the most promising from a crystallographic point of view. It would also be of interest to compare two so closely related compounds, one with sulfur and one with selenium as the central atom.

### Experimental

By cooling an alcoholic solution of trimethylselenonium iodide, crystals suitable for X-ray analysis were obtained in the form of colorless prisms, elongated in the direction of the  $a$  axis.

Although normally the substance is perfectly stable at room temperature, it was found that exposure to X-rays caused rapid decomposition unless the crystal was cooled below  $-15^\circ\text{C}$ .

For this reason only a set of comparatively poor  $h0l$  data was collected at room temperature.

All photographs used for determination of the cell dimensions and for three-dimensional intensity data were prepared with the crystals kept in the temperature

range between  $-15^\circ$  and  $-20^\circ\text{C}$ . The crystals were sealed in thin-walled glass capillaries.

Weissenberg and oscillation photographs indicated an orthorhombic unit cell. The following systematic extinctions were observed:  $0kl$  with  $k+l$  odd,  $hk0$  with  $h$  odd. The space group is accordingly  $Pnma$  or  $Pn2_1a$ . The cell dimensions were determined by means of zero level  $0kl$  and  $h0l$  Weissenberg photographs taken with unfiltered Cu radiation. Powder diagrams of KCl were superimposed on both ends of these films for calibration purposes. The  $2\theta$  values (based on  $a=6.2850$  Å for KCl at  $-15^\circ\text{C}$ ) for about 100 reflections were used in a least-squares refinement of the lattice parameters. The program used was written by R. A. Sparks at UCLA. With the wavelengths for Cu  $K$  radiation taken as  $\alpha_1=1.54050$ ,  $\alpha_2=1.54434$ , and  $\beta=1.39217$  Å, the following results were obtained:

$$a=14.078, b=8.000, c=6.177 \text{ \AA.}$$

Formal e.s.d.  $\sim 0.01\%$ .

The density measured by flotation was  $2.37 \text{ g.cm}^{-3}$ , while that calculated for  $Z=4$  is  $2.396 \text{ g.cm}^{-3}$ .

Integrated equi-inclination Weissenberg photographs of the layers with  $l=0$  through 6,  $k=0$  and 1, and  $h=0$  were prepared by use of filtered Mo radiation, with tin foil interleaves between the films. About two-thirds of the intensities recorded were measured with a densitometer, while the weakest one-third was estimated visually by use of a calibrated integrated comparison strip. Of about 580 unique reflections that were recorded, the majority was observed in four different octants. The intensities were corrected for Lorentz and polarization effects in the usual way.

The crystals used were roughly cut to approximate a spherical shape with a diameter of about 0.1–0.15 mm. No absorption corrections were applied.

### Solution and refinement of the structure

On the  $c$ -axis photographs it was noted that the intensity distribution was almost the same on all rows with

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$k$  even, and likewise for the rows with  $k$  odd. This observation suggested that the heavy atoms were located in the mirror planes at  $y = \frac{1}{4}$  and  $y = \frac{3}{4}$  in  $Pnma$ . The solution of the structure was based on this assumption. The successful refinement of the structure indicates that the choice of space group was correct.

With the  $h0l$  data collected at room temperature the Patterson function  $P(U, W)$  was calculated. The I–Se vector nearly parallel to the  $c$  axis, with the end point at  $U \sim 0.01$  and  $W \sim 0.33$  was readily identified. However, there appeared to be four possible solutions, and instead of trying them out one after one, a search program written by R. A. Sparks for the IBM 7090 was used to move the identified I–Se vector across one asymmetric unit. For each position the weighted sum of  $(F_o - F_c)^2$  was printed out. A scan of 2500 points required about the same amount of machine time as one two-dimensional Fourier synthesis. Four minima in the resulting table, corresponding to the four possible solutions of the Patterson function were observed. The I and Se positions derived from the deepest of these minima was used as input for a structure-factor Fourier-synthesis calculation. The resulting electron density projection clearly showed peaks corresponding to I, Se, one double C and one single C atom. The  $x, z$  parameters obtained were used as starting parameters for a two-dimensional least-squares refinement, which after a few cycles converged (with  $R = 0.067$ ) towards a set of positional parameters which were essentially the same as the final ones for I, Se and the 'double' C atom, whereas the C atom in the mirror plane showed a deviation of a couple of tenths of an Å.

Assuming a Se–C distance of 1.95 Å, and an angle C–Se–C of 100°, the  $y$  coordinate for the C atom not situated in a mirror plane was calculated.

Using the complete set of three-dimensional data, eight cycles of full-matrix least-squares refinement were performed. Following cycle three, anisotropic temperature factors for I and Se were introduced. It was

Table 1. *Final positional parameters and their standard deviations*

The e.s.d. (in parentheses) have been multiplied by 10<sup>4</sup>.

Atom	$x$	$y$	$z$
I	0.0922 (0.8)	0.75	0.4027 (2)
Se	0.0794 (1.2)	0.25	0.0672 (3)
C(1)	0.1705 (16)	0.25	0.3075 (37)
C(2)	0.1334 (10)	0.0649 (20)	–0.0961 (27)

Table 2. *Final thermal parameters (Å<sup>2</sup>) and their standard deviations*

The e.s.d. (in parentheses) have been multiplied by 10<sup>3</sup>. The relation between the  $B$ 's given below and the  $\beta$ 's in the anisotropic temperature factor expression is given by  $B_{11} = 4a^2\beta_{11}$ ,  $B_{13} = 2ac\beta_{13}$  etc.

Atom	$B_{11}$	$B_{22}$	$B_{33}$	$B_{13}$
I	2.497 (46)	2.747 (52)	2.773 (58)	–0.269 (43)
Se	2.018 (66)	3.796 (92)	2.201 (83)	0.271 (56)
C(1)	3.51 (420)	(isotropic)		
C(2)	3.40 (254)	(isotropic)		

not attempted to locate the hydrogen atoms, and they were not included in any calculations. Reflections too weak to be observed were not used in the least-squares refinement. The shifts in the last cycle were in no case greater than 1/20 of the corresponding standard deviation. The final parameters arrived at are given in Tables 1 and 2 together with their estimated standard deviations. The  $R$  index is 0.056 for all observed reflections. In Table 3 is given a comparison of observed and calculated structure factors. The  $F_c$ 's were calculated from the final parameters given in Tables 1 and 2. Unobserved reflections, indicated by  $U$ , are included in the table, with an  $F_o$  corresponding to the minimum observable intensity.

As a final check of the structure, a difference Fourier synthesis was calculated with all atoms subtracted. No disturbing features could be detected. The Fourier program used was one written by P. K. Gantzel and H. Hope at UCLA, and the least-squares program was ACA Computer Program No. 317 (UCLALS-1) written by P. K. Gantzel, R. A. Sparks and K. N. Trueblood; both programs were adapted for use on the UNIVAC 1107 computer. The least-squares program minimizes the weighted sum of the squares of  $KF_o - G(F_c)$  by a full-matrix routine. Of the scale factors  $K$  and  $G$ , the latter is one of the adjustable parameters. The weighting scheme of Hughes (1941) was used with  $4F_o(\min) = 58$ . The atomic form factors were those of Hanson, Herman, Lea & Skillman (1964). The isotropic temperature factors are of the form  $\exp(-B \sin^2 \theta/\lambda^2)$ , the anisotropic of the form  $\exp[-(h^2\beta_{11} + k^2\beta_{22} + l^2\beta_{33} + hk\beta_{12} + hl\beta_{13} + kl\beta_{23})]$ .

The  $R$  index is defined by  $R = \Sigma ||F_o| - |F_c|| / \Sigma |F_o|$ . The standard deviations were estimated from the inverse matrix of the normal equations.

### Description of the structure

In Fig. 1 are shown two projections of the structure, along the  $c$  axis and the  $b$  axis. Interatomic distances are given in Table 4 and Fig. 1(b). Angles are given in Table 4.

An electron diffraction study of dimethyl selenide (Goldish, Hedberg, Marsh & Shoemaker, 1955) gave a C–Se distance of 1.977 Å, and a C–Se–C angle of 98°. Comparison with the distances and angles found in the present study shows that no great change has taken place in these parameters with the formation of the selenonium ion. Within the limits of error the two crystallographically nonequivalent C–Se distances are equal, and they do not differ significantly from the sum of the single covalent bond radii (1.94 Å). The two C–Se–C angles are also almost equal, so that there is no indication of a difference in bonding between the two equatorial C atoms on one side, and the axial C atom (situated in the mirror plane) on the other. As in the trimethylsulfonium ion (Zuccaro & McCullough, 1959) the symmetry closely approximates  $3m$  if only the 'onium' ion is considered. However, if the position

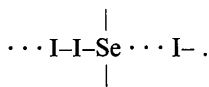


Table 4. *Interatomic distances and bond angles*

E.s.d. in parentheses	
Atoms	Distance
Se-C(1)	1.962 (0.023) Å
Se-C(2)	1.946 (0.016)
I-Se	3.776 (0.002)
I'-C(1)	4.173 (0.022)
I'-C(2)	4.137 (0.014)
Angle	
C(1)-Se-C(2)	97.9° (0.7°)
C(2)-Se-C(2')	99.1 (0.7)
C(1)-Se-I	178.9 (0.3)
C(2)-Se-I	81.5 (0.5)

of the iodine atom relative to the selenium group is taken into account it appears that this is not the true symmetry of the group. The iodine atom is situated in an axial position, at a distance of 3.78 Å from the Se atom. With a sum of van der Waals radii of 4.15 Å the observed interatomic distance, together with the atomic orientation, suggest the existence of a weak bond between selenium and iodine. It therefore seems reasonable to describe trimethylselenium iodide as a charge transfer complex, with the iodide ion acting as donor, and the selenium ion as acceptor.

It is interesting to compare this complex with the tetrahydroselenophene-iodine complex (Hope & McCullough, 1964) where a selenium-iodine bond of similar length (3.64 Å) and direction was found in the arrangement



There is also a resemblance between the present selenium salt and the compounds di-*p*-tolylselenium dichloride and di-*p*-tolylselenium dibromide (McCul-

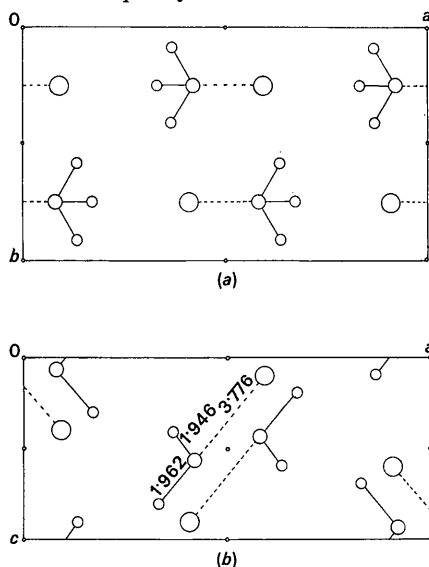


Fig. 1. (a) Projection of the structure of trimethylselenium iodide along the *c* axis. The broken lines indicate an assumed Se-I bond. In order of decreasing size, the circles represent I, Se and C. (b) Projection along the *b* axis. The crystallographically independent bond distances are given.

lough & Marsh, 1950*b*), where the halogen atoms occupy axial positions, while the organic groups occupy the equatorial positions. Apparently the selenium iodide possesses some of the geometrical characteristics of four-bonded selenium compounds.

Inspection of the thermal parameters reveals that the iodine atom vibrates very nearly isotropically, whereas the motion of the selenium atom is anisotropic with the largest component along the *b* axis. The reason for this, whether it indicates a slight disorder or not, is not known. Except for the Se-I distance, the packing distances are normal. An iodine atom rests against the three methyl groups, with one observed distance of 4.17 Å and two of 4.14 Å.

A comparison of the crystal structure of trimethylsulfonium iodide (Zuccaro & McCullough, 1959) with the present structure shows a striking similarity in the stacking of ions in the direction of the *b* axis [the central part of Fig. 1(*b*)]. It is interesting to note that even though the S-I distance of 3.89 Å is longer than the corresponding Se-I distance, it is still somewhat shorter than the sum of van der Waals radii (4.0 Å). With the similarity of the two structures in mind, it seems reasonable to assume that also in the trimethylsulfonium iodide we have a case of an ion-ion charge transfer complex.

Another example along the same lines is thiuret hydriodide. In the structure determination by Foss & Tjomsland (1958) an interionic S-I distance of only 3.62 Å was found. Hordvik & Sundsfjord (1965) have interpreted this as an indication of an ion-ion charge transfer bond.

The author is indebted to Professor J. D. McCullough for the gift of a sample of trimethylselenium iodide, for his interest and many inspiring discussions, to Mr Tore Onstad for help with the intensity measurements, to the National Science Foundation for financial support under Grant NSF-G24997, to the Fulbright Foundation for a travel grant, to Norges almenvitenskapelige forskningsråd for financial support under Grant D. 302-46, and to the Norwegian Computing Center for the use of the UNIVAC 1107 computer.

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